

# Geometric Quantization

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# You are used to quantizing

## Canonical Quantization

$$Q(q) = q$$

$$Q(p) = -i\hbar\partial_q$$

Nice properties:

- $[Q(q), Q(p)] = i\hbar \implies$  Uncertainty
- Standard potentials work nicely Ex:  $H = p^2/2m + V(q)$

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Bad properties:

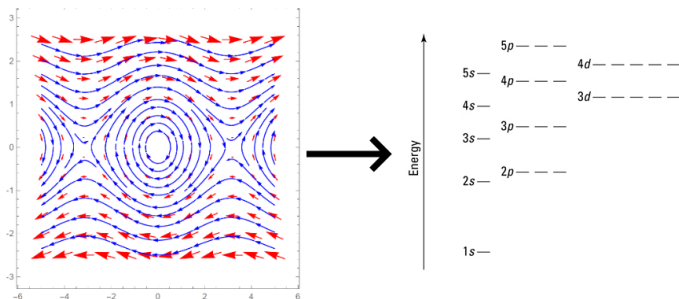
- $Q(qp)$  ambiguous
- Meanwhile, classical symmetries in  $C^\infty(M, \omega)$  w/ Poisson

# Quantizations should be nice

## Definition

A *quantization* is a map  $Q$  from smooth functions on a manifold to operators on a Hilbert Space obeying Dirac's axioms:

- $Q(1)=\text{Identity}$
- $[Q(A), Q(B)] = i\hbar Q(\{A, B\})$
- $Q$  of complete set of functions is a CSCO



# Quantizations are Hard

## Theorem (Groenewold-van Hove)

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Solution is to **throw something away**.

Deformation: let  $[Q(A), Q(B)] = i\hbar Q(\{A, B\}) + O(\hbar^2)$  - fine for chemists but Ehrenfest would be unhappy.

Geometric: make oversized Hilbert space then throw away pieces (like restricting to  $\{f(q)p + g(q)\}$ )

## More Naive Quantizations

The “I just learned symplectic geometry” quantization

Know  $\{A, B\} = \omega(v_A, v_B)$ .

Try  $Q(f) = v_f$ .

Then  $Q(\{A, B\}) = v_{\{A, B\}} = [v_A, v_B] = [Q(A), Q(B)]$

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The “uncertain” quantization

Try  $Q(f) = v_f + f$ .

But  $[Q(q), Q(p)] = 2 \neq 1 = Q(1)$



# Geometric Extensions: Forms and Curvature

## The Hove-Segal quantization

Know, locally,  $\omega = d\theta$

Try  $Q(f) = v_f + f \cdot + \theta(v_f) \cdot$

Works on Hilbert Space  $\mathbb{R}^{2n} \rightarrow \mathbb{C}$ ,  $(f, g) = \int_M f^* g \omega^n$

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## The Kirillov-Konstant-Soriau quantization on $(M, \omega)$

Take prequantum line bundle  $U(1) \rightarrow Q \xrightarrow{\pi} M$

$$Q(f) = -i\hbar \nabla_{v_f} + f$$

With  $\nabla_{v_f} = (d + A)(v_f)$  and curvature  $F = dA$ ,  
if  $F = \omega$ , then

$$[-i\nabla_A + A, -i\nabla_B + B] = \dots = -i(-i\nabla_{\{A,B\}} + \{A, B\})$$

## Quantumness and Angular Momentum

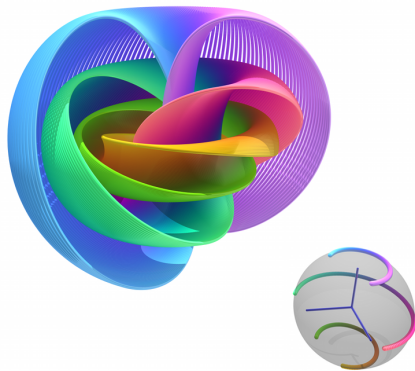
Note that  $F$  is  $U(1)$  curvature iff  $\int_{\alpha} F \in 2\pi\mathbb{Z}, \forall \alpha \in H_2(M)$ .

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Consider space  $U(1) \longrightarrow S^3 \xrightarrow{\langle \vec{\sigma} \rangle} S^2$  with Berry Connection

$$A = -is(\langle \partial_{\theta} \rangle d\theta + \langle \partial_{\phi} \rangle d\phi) \quad \omega = \frac{s}{2} \sin(\theta) d\theta \wedge d\phi \quad \int_{S^2} \omega = 2\pi s$$

Thus spin quantized.

Recover angular momentum algebra from  $-s/2\{x, y, z\}$

Hilbert space:  $\frac{1}{f(z)}(z)^{2s}$  is holomorphic  $\implies f \in \text{span}\{1, z, \dots, z^{2s}\}$ ,  
 $z = \tan(\theta/2)e^{i\phi}$

Thank you.  
Any Questions?

## Messy Ideas

The nonexample:

$$Q(\{q^3, p^2\} + \frac{1}{12}\{\{p^2, q^3\}, \{q^2, p^3\}\}) \neq [q^3, p^2] + \frac{1}{12}[[p^2, q^3], [q^2, p^3]]$$

Angular Momenta:

$$Q(x) = \frac{-s}{2} \cos \phi \tan(\theta/2) + i\hbar \cos \phi \cot \theta \partial_\phi - i\hbar \sin \phi \partial_\theta$$

$$Q(y) = \frac{-s}{2} \sin \phi \tan(\theta/2) + i\hbar \sin \phi \cot \theta \partial_\phi - i\hbar \cos \phi \partial_\theta$$

$$Q(z) = -i\hbar \partial_\phi - \frac{s}{2}$$

Completeness: Exists a dense set of separately analytic wavefunctions for quantization of irreducible base of Poisson algebra ( $\sum |Q(b)^k \psi| t^k / k!$  converges in  $\mathcal{H}$  for some  $t$  for each  $b$ ).

## Can also quantize states directly

Use WKB approximation  $\psi \sim A(q)e^{iS(q)/\hbar}$

Evolution via transport requires half-density  $A(q)$

Projectible on Lagrangian immersion  $\omega_L = 0$  with asymptotic inverse Fourier transform.





Polarizations have issues on overlaps so need Maslov correction.

$$\psi(Q, t) = \sum_{\{q_j: g^t q_j = Q\}} \varphi(q_j) \left| \frac{D[q(g^t(p(q_j)), q_j)]}{Dq_j} \right|^{-1/2} \exp \left\{ i/\hbar \left[ s(q_j) + \int_0^t L(g^t(p(q_i)), q_i) dt \right] - \frac{i\pi\mu_j}{2} \right\}$$

Where  $S(q) = (Sq, q)/2$  and  $\mu_j$  sign flips in  $(0, t]$  of  $\det(I - iS)^{-1}(I + iS)$  along  $(-H(g^t(q)), g^t(p), t, g^t(q))$ .



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